

Page 187, equations (4.208) and (4.209)

From the definition (4.165) of coefficients $\omega_{\beta\gamma}$ it follows:

$$\begin{aligned} \sum_{\gamma=1}^n \omega_{\beta\gamma} d\rho_{\gamma} &= \sum_{\gamma=1}^n \left(g_{\gamma} \delta_{\beta\gamma} - \rho_{\beta} \frac{\partial \hat{f}_{\beta}}{\partial \rho_{\gamma}} - f_{\beta} \delta_{\beta\gamma} \right) d\rho_{\gamma} \\ &= g_{\beta} d\rho_{\beta} - \rho_{\beta} \sum_{\gamma=1}^n \frac{\partial \hat{f}_{\beta}}{\partial \rho_{\gamma}} d\rho_{\gamma} - f_{\beta} d\rho_{\beta}; \quad \beta = 1, \dots, n-1. \end{aligned} \quad (1)$$

The differential of partial free energy reads:

$$df_{\beta} = \frac{\partial \hat{f}_{\beta}}{\partial T} dT + \sum_{\gamma=1}^n \frac{\partial \hat{f}_{\beta}}{\partial \rho_{\gamma}} d\rho_{\gamma}$$

and from this then follows:

$$-\rho_{\beta} \sum_{\gamma=1}^n \frac{\partial \hat{f}_{\beta}}{\partial \rho_{\gamma}} d\rho_{\gamma} = \rho_{\beta} \frac{\partial \hat{f}_{\beta}}{\partial T} - \rho_{\beta} df_{\beta}. \quad (2)$$

Additional relationships among differential can be obtained from the relationship (4.193) for partial thermodynamic pressure:

$$\begin{aligned} dP_{\beta} &= \rho_{\beta} dg_{\beta} + g_{\beta} d\rho_{\beta} - \rho_{\beta} df_{\beta} - f_{\beta} d\rho_{\beta}, \\ g_{\beta} d\rho_{\beta} - \rho_{\beta} df_{\beta} - f_{\beta} d\rho_{\beta} &= dP_{\beta} - \rho_{\beta} dg_{\beta}. \end{aligned} \quad (3)$$

Substituting from (2) and (3) into (1) we obtain:

$$\begin{aligned} \sum_{\gamma=1}^n \omega_{\beta\gamma} d\rho_{\gamma} &= g_{\beta} d\rho_{\beta} + \rho_{\beta} \frac{\partial \hat{f}_{\beta}}{\partial T} dT - \rho_{\beta} df_{\beta} - f_{\beta} d\rho_{\beta} \\ &= dP_{\beta} - \rho_{\beta} dg_{\beta} + \rho_{\beta} \frac{\partial \hat{f}_{\beta}}{\partial T} dT \end{aligned} \quad (4)$$

and this is equation (4.208).

Summing (4) from $\beta = 1$ to $\beta = n-1$ we get:

$$\sum_{\beta=1}^{n-1} \sum_{\gamma=1}^n \omega_{\beta\gamma} d\rho_{\gamma} = \sum_{\beta=1}^{n-1} dP_{\beta} - \sum_{\beta=1}^{n-1} \rho_{\beta} dg_{\beta} + \sum_{\beta=1}^{n-1} \rho_{\beta} \frac{\partial \hat{f}_{\beta}}{\partial T} dT. \quad (5)$$

From eq. (4.187) we have

$$dP = \sum_{\alpha=1}^n dP_{\alpha} = \sum_{\beta=1}^{n-1} dP_{\beta} + dP_n. \quad (6)$$

From Gibbs-Duhem equation (4.207), taking into account the definitions of mass fraction, (4.22), and volume, (4.195), it further follows:

$$\begin{aligned} -\rho s dT + \rho v dP - \sum_{\alpha=1}^n \rho_{\alpha} dg_{\alpha} &= -\rho s dT + dP - \sum_{\beta=1}^{n-1} \rho_{\beta} dg_{\beta} - \rho_n dg_n \\ &= 0. \end{aligned} \quad (7)$$

Substitution from (6) and (7) into (5) results in following expression:

$$\sum_{\beta=1}^{n-1} \sum_{\gamma=1}^n \omega_{\beta\gamma} d\rho_{\gamma} = dP - dP_n + \rho_n dg_n + \rho s dT - dP + \sum_{\beta=1}^{n-1} \rho_{\beta} \frac{\partial \hat{f}_{\beta}}{\partial T} dT. \quad (8)$$

From equation (4.164), taking into account (4.92) and the definition of mass fraction, (4.22), we derive:

$$\begin{aligned} \rho s dT &= -\rho \frac{\partial \hat{f}}{\partial T} dT = -\rho \frac{\partial}{\partial T} \left(\sum_{\alpha=1}^n w_{\alpha} \hat{f}_{\alpha} \right) dT = -\rho \left(\sum_{\alpha=1}^n w_{\alpha} \frac{\partial \hat{f}_{\alpha}}{\partial T} \right) dT \\ &= -\sum_{\alpha=1}^n \rho_{\alpha} \frac{\partial \hat{f}_{\alpha}}{\partial T} dT. \end{aligned} \quad (9)$$

Substituting (9) into (8), eq. (4.209) follows:

$$\begin{aligned} \sum_{\beta=1}^{n-1} \sum_{\gamma=1}^n \omega_{\beta\gamma} d\rho_{\gamma} &= -dP_n + \rho_n dg_n - \sum_{\alpha=1}^n \rho_{\alpha} \frac{\partial \hat{f}_{\alpha}}{\partial T} dT + \sum_{\beta=1}^{n-1} \rho_{\beta} \frac{\partial \hat{f}_{\beta}}{\partial T} dT \\ &= -dP_n + \rho_n dg_n - \rho_n \frac{\partial \hat{f}_n}{\partial T} dT. \end{aligned}$$